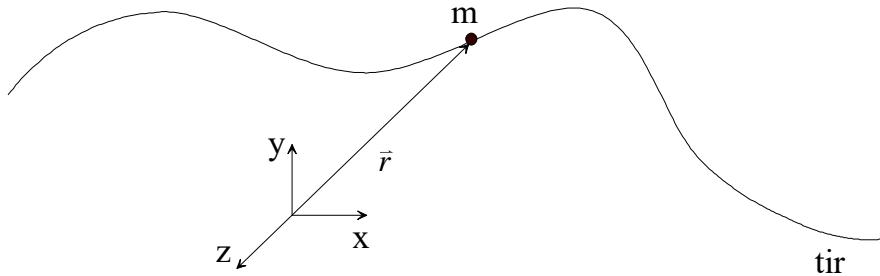


Izrek o vrtilni količini

Točkasto telo

$$\vec{r}(t) = (x, y, z)$$



$$\text{hitrost } \vec{v} = \frac{d\vec{r}}{dt} \quad \text{pospešek } \vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2\vec{r}}{dt^2}$$

II. Newton-ov zakon za gibanje točkastega telesa:

$$m\vec{a} = \vec{F}$$

$$\vec{r} \times m\vec{a} = \vec{r} \times \vec{F} = \vec{M}$$

\vec{M} = **rezultantni navor vseh zunanjih sil**

$$\frac{d}{dt}(\vec{r} \times m\vec{v}) = \frac{d\vec{r}}{dt} \times m\vec{v} + \vec{r} \times m \frac{d\vec{v}}{dt} = \vec{v} \times m\vec{v} + \vec{r} \times m\vec{a} = \vec{r} \times m\vec{a}$$

$$\frac{d(\vec{r} \times m\vec{v})}{dt} = \vec{M}$$

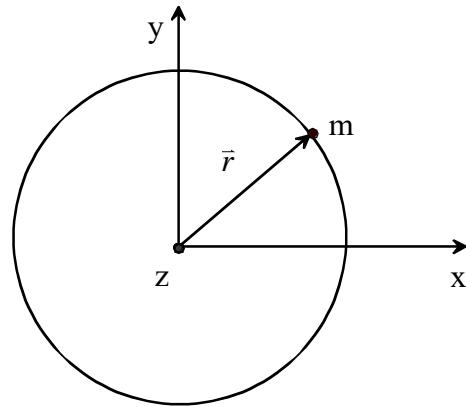
vrtilna količina točkastega telesa: $\vec{\Gamma} = \vec{r} \times m\vec{v}$

$$\frac{d\vec{\Gamma}}{dt} = \vec{M}$$

$$d\vec{\Gamma} = \vec{M} dt$$

$$\int_{t_1}^{t_2} \vec{M} dt = \vec{\Gamma}_2 - \vec{\Gamma}_1 = \Delta \vec{\Gamma}$$

Primer: enakomerno **krožečo točkasto maso**



$$\vec{\Gamma} = \vec{r} \times m \vec{v} \quad v = \omega r$$

$$\Gamma = rmv = rm\omega r = mr^2\omega \quad \vec{\Gamma} = J \vec{\omega}$$

$$W_k = \frac{1}{2}mv^2 = \frac{1}{2}mr^2\omega^2 = \frac{1}{2}J\omega^2$$

kjer definiramo **vztrajnostni moment krožeče točkaste mase**

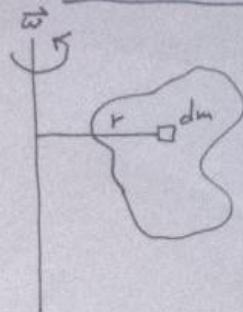
$$J = mr^2$$

$$\vec{M} = \frac{d\vec{\Gamma}}{dt} = J \frac{d\vec{\omega}}{dt} = J \vec{\alpha} \quad \vec{\alpha} = \frac{d\vec{\omega}}{dt}$$

$$\int_{t_1}^{t_2} \vec{M} dt = J \vec{\omega}_2 - J \vec{\omega}_1$$

Kinetična energija točke telesa,

ki se vrti s kotno hitrostjo deski filzne osi



za vsak delik z maso dm velja:

$$V = r\omega$$

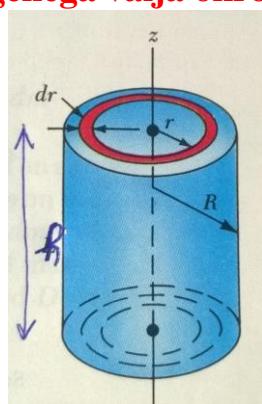
$$W_k = \frac{1}{2} \int v^2 dm = \frac{1}{2} \int (r\omega)^2 dm = \frac{1}{2} \omega^2 \underbrace{\int r^2 dm}_{J} = \\ = \frac{1}{2} \omega^2 J$$

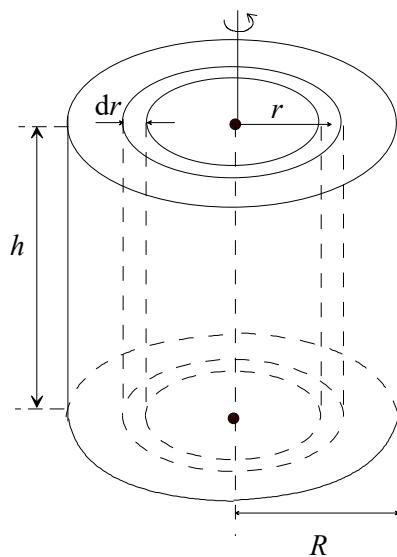
$J = \int r^2 dm$ = vztrajnostni moment
glede na izbrano (fiksno) os

$$W_k = \frac{1}{2} J \omega^2$$

velja ne glede na izbrano os
Os je lahka izven telesa, tj. roči

Zgled: vztrajnostni moment **homogenega valja okrog geometrijske osi**





$R \equiv$ polmer valja

$h \equiv$ višina valja

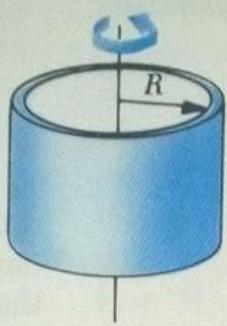
$$J = \int r^2 dm = \int r^2 \rho dV \quad \rho = \text{gostota valja}$$

$$dV = 2\pi r dr h$$

$$J = \int_0^R r^2 \rho 2\pi r h dr = 2\pi h \rho \int_0^R r^3 dr = 2\pi h \rho \frac{r^4}{4} \Big|_0^R = \pi R^4 \rho \frac{h}{2} = \frac{1}{2} m R^2$$

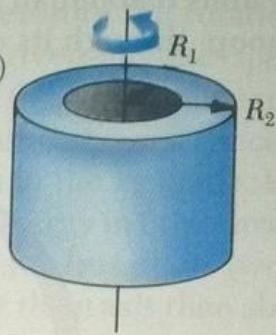
$$\text{celotna masa valja } m = \pi R^2 h \rho$$

Hoop or
cylindrical shell
 $I_c = MR^2$

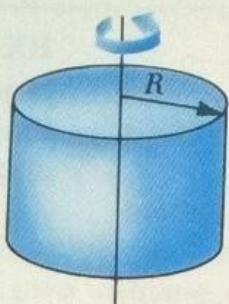


Hollow cylinder

$$I_c = \frac{1}{2} M(R_1^2 + R_2^2)$$

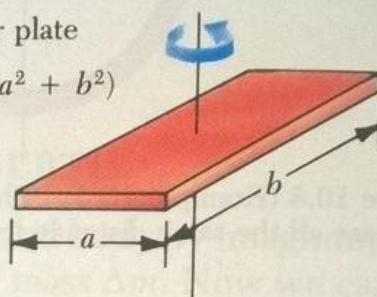


Solid cylinder
or disk
 $I_c = \frac{1}{2} MR^2$

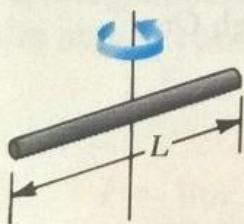


Rectangular plate

$$I_c = \frac{1}{12} M(a^2 + b^2)$$

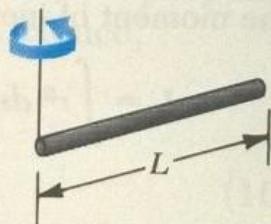


Long thin rod
 $I_c = \frac{1}{12} ML^2$

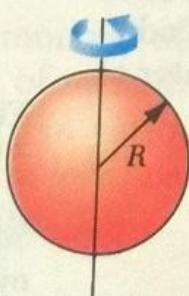


Long thin rod

$$I = \frac{1}{3} ML^2$$

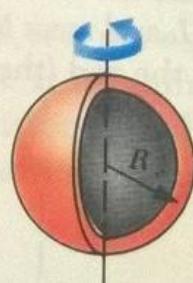


Solid sphere
 $I_c = \frac{2}{5} MR^2$



Thin spherical
shell

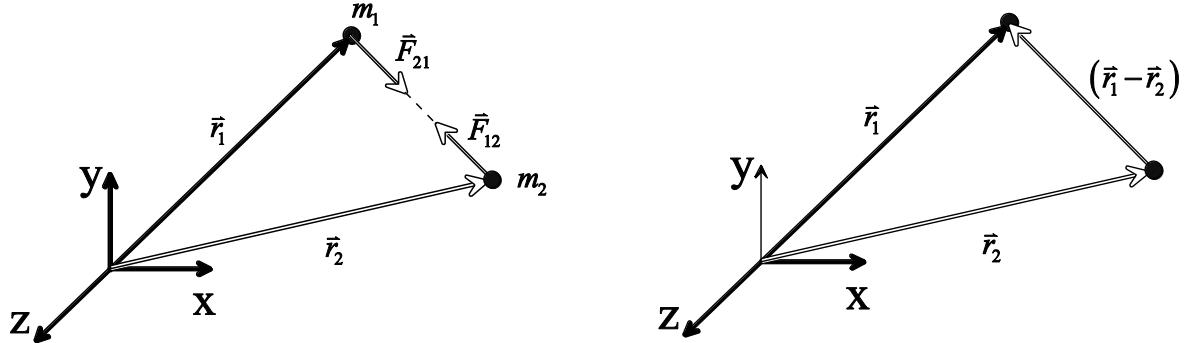
$$I_c = \frac{2}{3} MR^2$$



Sistem točkastih mas

Poseben primer : sistem dveh točkastih delcev

III. Newtonov zakon : $\vec{F}_{12} + \vec{F}_{21} = \vec{0}$



$$\sum_i \vec{M}_i = \vec{r}_1 \times \vec{F}_{21} + \vec{r}_2 \times \vec{F}_{12} = \vec{r}_1 \times \vec{F}_{21} + \vec{r}_2 \times (-\vec{F}_{21})$$

$$\sum_i \vec{M}_i = (\vec{r}_1 - \vec{r}_2) \times \vec{F}_{21} = \vec{0}$$

POSPLOŠITEV: navor vseh notranjih sil nič, če so dvodelčne notranje sile centralne

vrtilna količina množice točkastih mas :

$$\vec{\Gamma} = \sum_i \vec{r}_i \times m_i \vec{v}_i$$

$$\begin{aligned} \frac{d\vec{\Gamma}}{dt} &= \frac{d}{dt} \left(\sum_i \vec{r}_i \times m_i \vec{v}_i \right) = \sum_i \frac{d\vec{r}_i}{dt} \times m_i \vec{v}_i + \sum_i \vec{r}_i \times m_i \frac{d\vec{v}_i}{dt} = \\ &= \sum_i \vec{v}_i \times m_i \vec{v}_i + \sum_i \vec{r}_i \times m_i \vec{a}_i = \sum_i \vec{r}_i \times m_i \vec{a}_i \end{aligned}$$

II. Newtonov zakon za posamezno točkasto maso :

$$m_i \vec{a}_i = \sum_j \vec{F}_{ji} + \vec{F}_i$$

$\sum_j \vec{F}_{ji}$ vsota vseh notranjih sil, ki delujejo na maso m_i
 \vec{F}_i pa rezultanta vseh zunanjih sil, ki delujejo na maso m_i

$$\frac{d\vec{\Gamma}}{dt} = \sum_i \vec{r}_i \times m_i \vec{a}_i \quad m_i \vec{a}_i = \sum_j \vec{F}_{ji} + \vec{F}_i$$

$$\frac{d\vec{\Gamma}}{dt} = \sum_i \left(\vec{r}_i \times \sum_j \vec{F}_{ji} \right) + \sum_i \left(\vec{r}_i \times \vec{F}_i \right)$$

KER vsota notranjih navorov $\sum_i \left(\vec{r}_i \times \sum_j \vec{F}_{ji} \right) = 0$

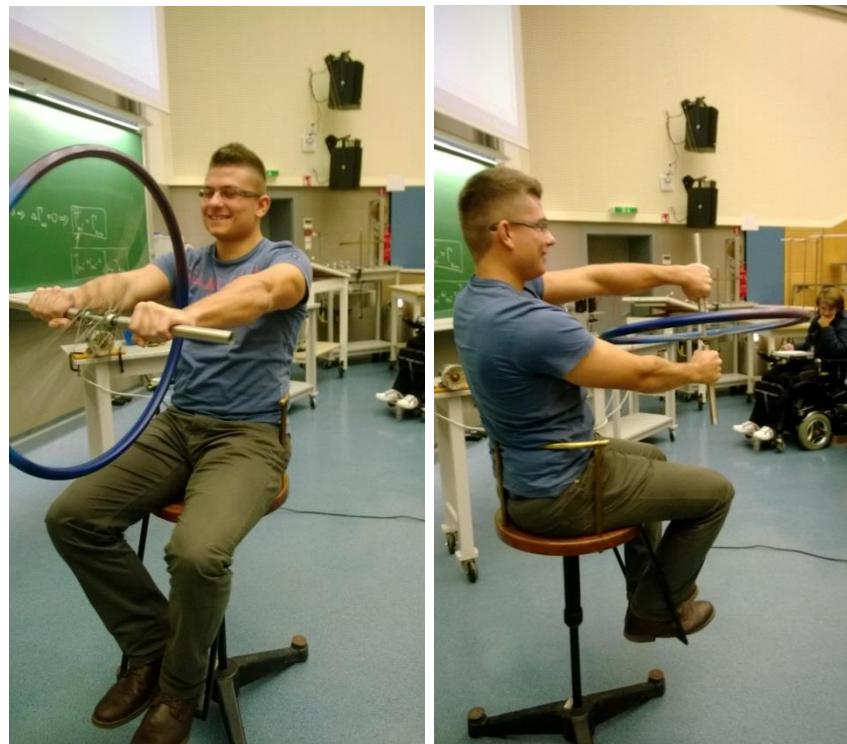
$$\frac{d\vec{\Gamma}}{dt} = \sum_i \left(\vec{r}_i \times \vec{F}_i \right) = \sum_i \vec{M}_i$$

izrek o vrtilni količini za sistem točkastih mas

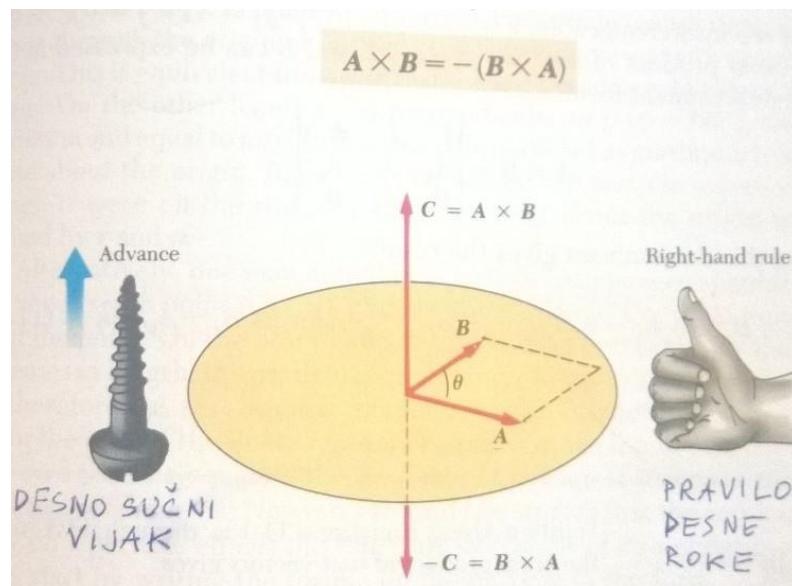
$$\vec{\Gamma} = \sum_i \vec{r}_i \times m_i \vec{v}_i$$

$\sum_i \vec{M}_i$ = vsota vseh zunanjih navorov, ki delujejo na sistem
točkastih mas m_1, m_2, m_3, \dots

ZGLED: človek na vrtljivem stolu, z rokami drži kolo od bicikla:



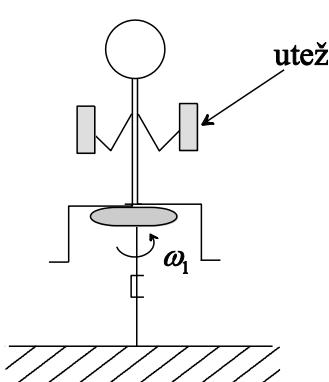
$$\vec{\Gamma} = \sum_i \vec{r}_i \times m_i \vec{v}_i \quad \int \vec{M} dt = \Delta \vec{\Gamma} = 0 \quad \vec{\Gamma} = \text{konst.}$$



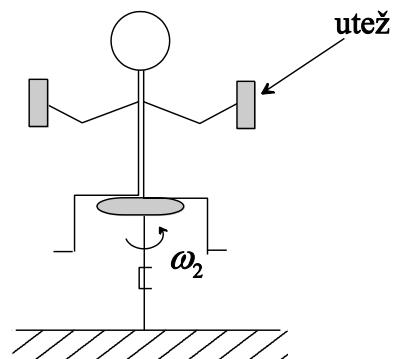
Zgled: človek na vrtljivem stolu z utežmi v obeh rokah (navor notranjih sil = 0)



$$\int_{t_1}^{t_2} \vec{M} \, dt = 0 \quad \int_{t_1}^{t_2} \vec{M} \, dt = J \vec{\omega}_2 - J \vec{\omega}_1 = 0 \quad J \vec{\omega}_2 = J \vec{\omega}_1$$



$$J_1 < J_2$$



vztrajnostni moment: J_1

vrtilna količina:

$$\Gamma_1 = J_1 \omega_1$$

vztrajnostni moment: J_2

vrtilna količina: $\Gamma_2 = J_2 \omega_2$

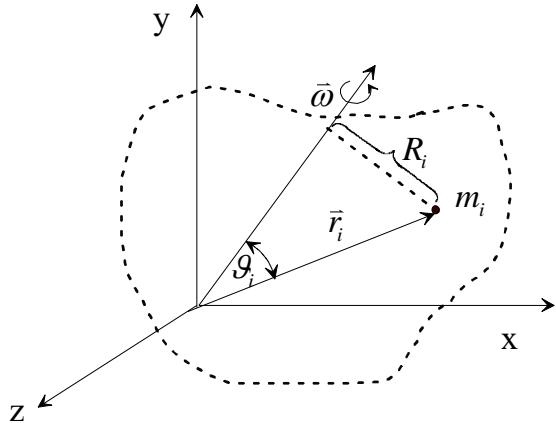
Navor notranjih sil enak nič:

$$J_1 \omega_1 = J_2 \omega_2$$

$$\omega_2 = \omega_1 \frac{J_1}{J_2} < \omega_1 \quad \text{če je : } J_1 < J_2$$

Togo telo

- $\bar{\omega}$ je trenutna kotna hitrost v smeri osi okoli katere se vrtijo vsi deli telesa (vse točkaste mase) z enako kotno hitrostjo: $\bar{\omega}(\vec{r}_i)=\omega=\text{konst.}$
- koordinatno izhodišče je na osi vrtenja



$$\bar{\omega} = (\omega_x, \omega_y, \omega_z)$$

$$\vec{v}_i = \bar{\omega} \times \vec{r}_i$$

$$\vec{r}_i = (x_i, y_i, z_i)$$

$$|\vec{v}_i| = \omega \underbrace{r_i}_{R_i} \sin \vartheta_i = \omega R_i$$

$$r_i^2 = x_i^2 + y_i^2 + z_i^2$$

sistem točkastih mas

$$\vec{\Gamma} = \sum_i \vec{r}_i \times m_i \vec{v}_i$$

$$\frac{d\vec{\Gamma}}{dt} = \sum_i \vec{M}_i \equiv \vec{M}$$

togo telo: $\sum_i \rightarrow \int , m_i \rightarrow dm , \vec{r}_i \rightarrow \vec{r} , \vec{v}_i \rightarrow \vec{v}$

$$\vec{\Gamma} = \int \vec{r} \times m_i \vec{v}_i \rightarrow \int \vec{r} \times \vec{v} dm$$

$$\frac{d\vec{\Gamma}}{dt} = \vec{M}$$

$$\vec{r} \times \vec{v} = \vec{r} \times (\bar{\omega} \times \vec{r}) = \bar{\omega}(\vec{r} \cdot \vec{r}) - \vec{r}(\vec{r} \cdot \bar{\omega}) , \text{ od koder sledi:}$$

$$\vec{\Gamma} = \int \vec{r} \times \vec{v} dm = \int \vec{r} \times (\bar{\omega} \times \vec{r}) dm = \int \bar{\omega} r^2 dm - \int \vec{r} (\bar{\omega} \cdot \vec{r}) dm$$

$$\bar{\Gamma} = \int \bar{r} \times \bar{v} \, dm = \int \bar{r} \times (\bar{\omega} \times \bar{r}) \, dm = \int \bar{\omega} r^2 \, dm - \int \bar{r} (\bar{\omega} \cdot \bar{r}) \, dm$$

$$\begin{aligned}\bar{r} &= (x, y, z) & r^2 &= \bar{r} \cdot \bar{r} = x^2 + y^2 + z^2 \\ \bar{\omega} &= (\omega_x, \omega_y, \omega_z) & \bar{\omega} \cdot \bar{r} &= \omega_x x + \omega_y y + \omega_z z\end{aligned}$$

zapis po komponentah:

$$\begin{aligned}\Gamma_x &= \omega_x \int r^2 \, dm - \int x (\omega_x x + \omega_y y + \omega_z z) \, dm \\ \Gamma_y &= \omega_y \int r^2 \, dm - \int y (\omega_x x + \omega_y y + \omega_z z) \, dm \\ \Gamma_z &= \omega_z \int r^2 \, dm - \int z (\omega_x x + \omega_y y + \omega_z z) \, dm\end{aligned}$$

preureditev:

$$\begin{aligned}\Gamma_x &= \omega_x \underbrace{\int (y^2 + z^2) \, dm}_{J_{xx}} - \omega_y \underbrace{\int xy \, dm}_{D_{xy}} - \omega_z \underbrace{\int xz \, dm}_{D_{xz}} \\ \Gamma_y &= -\omega_x \underbrace{\int xy \, dm}_{D_{yx}} + \omega_y \underbrace{\int (x^2 + z^2) \, dm}_{J_{yy}} - \omega_z \underbrace{\int zy \, dm}_{D_{yz}} \\ \Gamma_z &= -\omega_x \underbrace{\int zx \, dm}_{D_{zx}} - \omega_y \underbrace{\int zy \, dm}_{D_{zy}} + \omega_z \underbrace{\int (x^2 + y^2) \, dm}_{J_{zz}}\end{aligned}$$

vektorski zapis (vektorja $\bar{\Gamma}$ in $\bar{\omega}$ v splošnem nista vzporedna):

$$\bar{\Gamma} = \underline{J} \bar{\omega}$$

\underline{J} = vztrajnostni tenzor :

$$\underline{J} = \begin{bmatrix} J_{xx} & -D_{xy} & -D_{xz} \\ -D_{yx} & J_{yy} & -D_{yz} \\ -D_{zx} & -D_{zy} & J_{zz} \end{bmatrix}$$

J_{xx}, J_{yy}, J_{zz} so vztrajnostni momenti okrog treh pravokotnih osi
 $D_{xy}, D_{xz}, D_{yz}, D_{yx}, D_{zx}, D_{zy}$ so deviacijski momenti

$$J_{xx} = \int (y^2 + z^2) dm,$$

$$J_{yy} = \int (x^2 + z^2) dm,$$

$$J_{zz} = \int (x^2 + y^2) dm,$$

$$D_{xy} = \int xy dm, \quad D_{xz} = \int xz dm$$

$$D_{yx} = \int yx dm, \quad D_{yz} = \int yz dm$$

$$D_{zx} = \int zx dm, \quad D_{zy} = \int yz dm$$

$$D_{xy} = D_{yx}$$

$$D_{xz} = D_{zx}$$

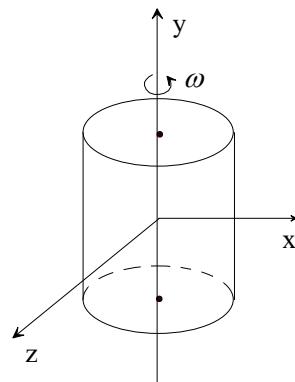
$$D_{yz} = D_{zy}$$

Diagonalizacija tenzorja \underline{J} : $D_{ij} = 0$, koordinate osi so usmerjene vzdolž glavnih osi:

$$\underline{J} = \begin{vmatrix} J_x & 0 & 0 \\ 0 & J_y & 0 \\ 0 & 0 & J_z \end{vmatrix}, \quad J_x \equiv J_{xx}, J_y \equiv J_{yy}, J_z \equiv J_{zz}$$

Poseben primer: $\bar{\omega} = (0, \omega, 0)$ v sistemu glavnih osi (\underline{J} je diagonaliziran)

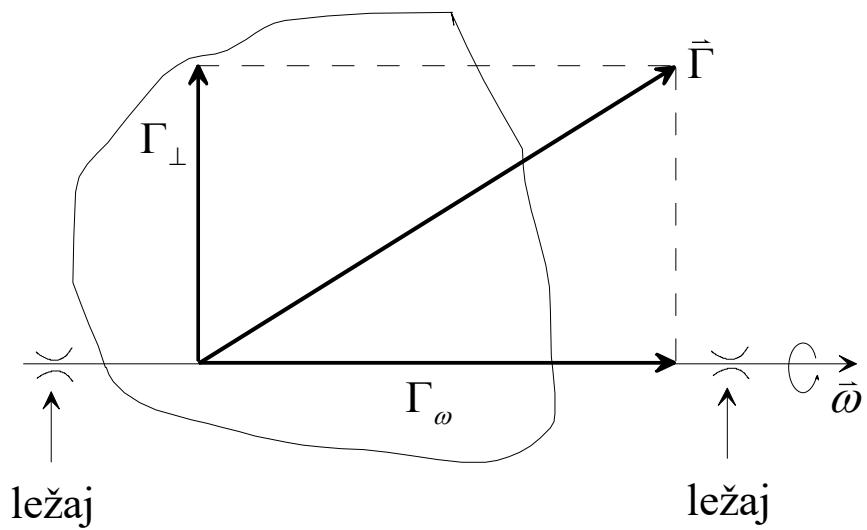
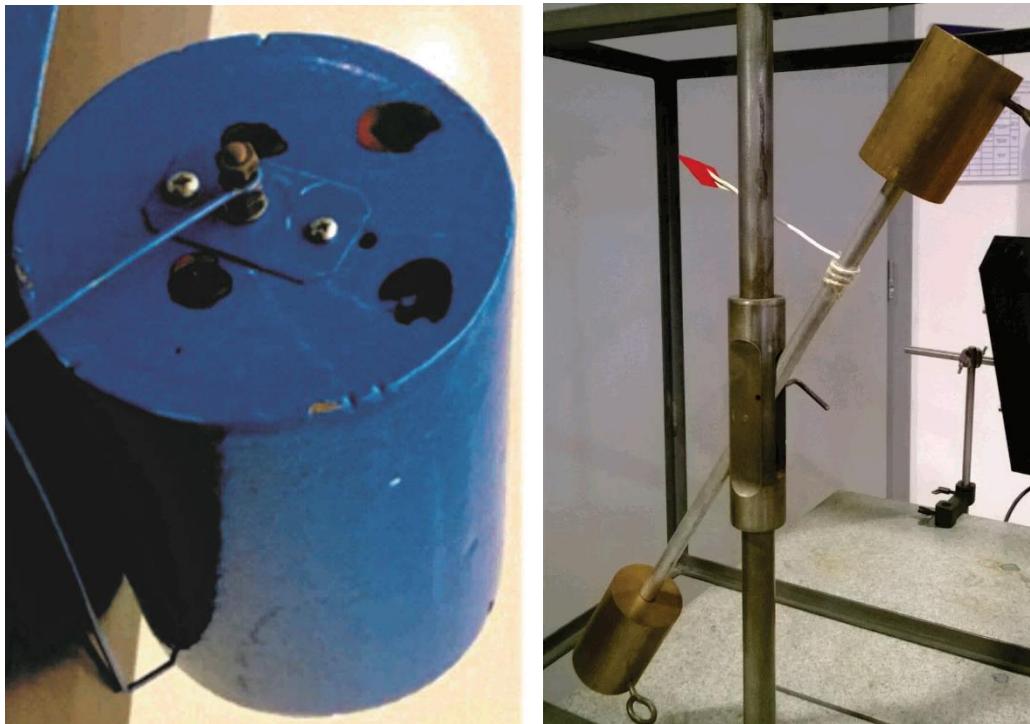
$\bar{\Gamma}$ in $\bar{\omega}$ sta vzporedna

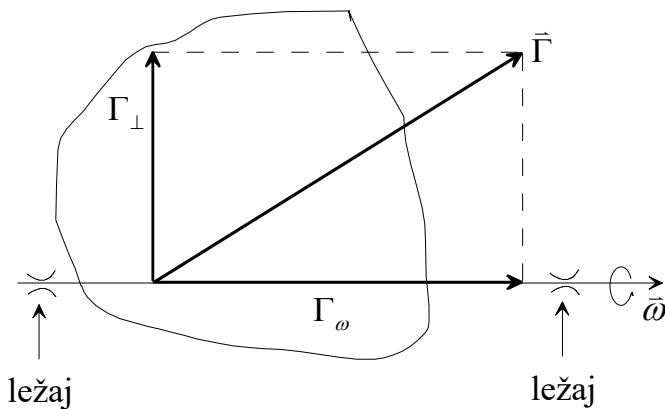


$$\Gamma_y = J_y \omega,$$

$$\frac{d\Gamma_y}{dt} = M_y \quad J_y = \int (x^2 + z^2) dm$$

Poseben primer: os vrtenja togega telesa je vpeta v ležaje





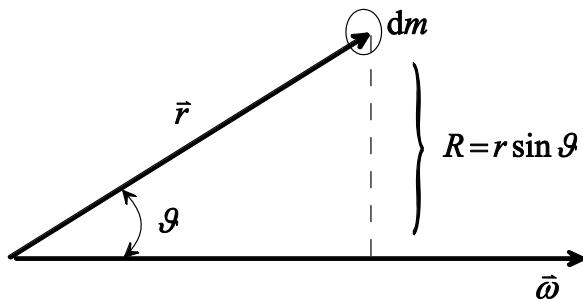
Komponenta vektorja $\vec{\Gamma}$ v smeri kotne hitrosti $\vec{\omega}$

$$\Gamma_\omega = \vec{\Gamma} \cdot \frac{\vec{\omega}}{\omega} \quad \left(\left| \frac{\vec{\omega}}{\omega} \right| = 1 \right)$$

M_ω = komponento navora \vec{M} v smeri vektorja kotne $\vec{\omega}$:

$$\begin{aligned} \Gamma_\omega &= \frac{\vec{\omega}}{\omega} \cdot \int \vec{r} \times \vec{v} \, dm = \frac{\vec{\omega}}{\omega} \cdot \int \vec{r} \times (\vec{\omega} \times \vec{r}) \, dm = \frac{\vec{\omega}}{\omega} \cdot \left[\int r^2 \vec{\omega} \, dm - \int \vec{r} (\vec{r} \cdot \vec{\omega}) \, dm \right] = \\ &= \int \left(r^2 \omega - \frac{(\vec{r} \cdot \vec{\omega})^2}{\omega} \right) \, dm = \int r^2 \omega (1 - \cos^2 \vartheta) \, dm = \int r^2 \omega \sin^2 \vartheta \, dm = \omega \int R^2 \, dm = \omega J_\omega \end{aligned}$$

$$\vec{r} \cdot \vec{\omega} = r \omega \cos \vartheta \quad R = r \sin \vartheta \quad J_\omega = \int R^2 \, dm$$



$$\Gamma_\omega = J_\omega \omega \quad M_\omega = \frac{d\Gamma_\omega}{dt} = J_\omega \frac{d\omega}{dt} = J_\omega \alpha$$

$$M_\omega = J_\omega \alpha \quad \alpha = \frac{d\omega}{dt} = \text{kotni pospešek}$$

$$M_\omega = \frac{d\Gamma_\omega}{dt} \quad \int M_\omega dt = \int_{\Gamma_1}^{\Gamma_2} d\Gamma_\omega = \Gamma_2 - \Gamma_1 = \Delta\Gamma$$

$$\int M_\omega dt = \Delta\Gamma = 0 \quad \Gamma_2 = \Gamma_1$$

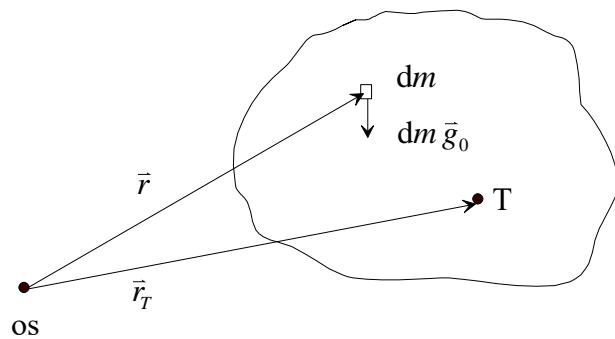
Poskus: tresenje v ležajih



Zgled: vreteno in utež



Navor sile teže telesa

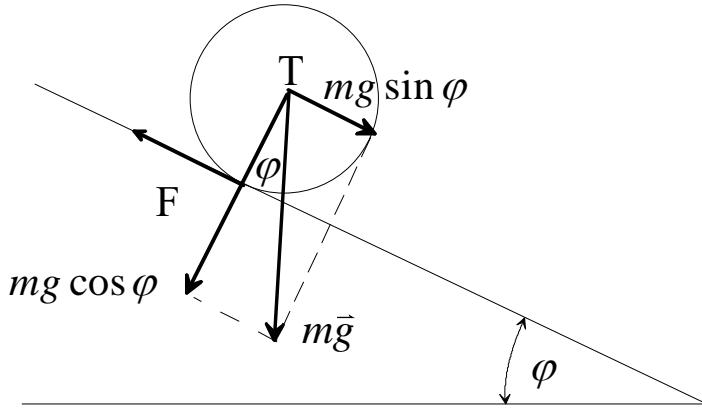


$T \equiv$ težišče telesa
 $dm \equiv$ masa majhnega koščka
 togega telesa

$$\begin{aligned}
 \vec{M}_g &= \int \vec{r} \times d\vec{F}_g = \int \vec{r} \times dm \vec{g}_0 = \int \vec{r} dm \times \vec{g}_0 = \left(\int \vec{r} dm \right) \times \vec{g}_0 = \\
 &= m \frac{\int \vec{r} dm}{m} \times \vec{g}_0 = m \vec{r}_T \times \vec{g}_0 = \vec{r}_T \times m \vec{g}_0 = \vec{r}_T \times \vec{F}_g
 \end{aligned}$$

$$\vec{r}_T = \frac{1}{m} \int \vec{r} dm \quad \vec{F}_g = \int dm \vec{g}_0 = m \vec{g}_0$$

Zgled: kotaljenje **homogenega valja** po klancu navzdol **brez podrsavanja**



$$mg \sin \varphi - F = m a_T$$

$$FR = J \alpha = \frac{1}{2} m R^2 \frac{a_T}{R} = \frac{1}{2} R m a_T$$

kotaljenje brez podrsavanja je hitrost težišča valja

$$J = \frac{1}{2} m R^2 \quad v_T = R \omega \quad a_T = R \alpha$$

$$mg \sin \varphi - \frac{1}{2} m a_T = m a_T \quad a_T = \frac{2}{3} g \sin \varphi$$

$$F = \frac{1}{2} m a_T = \frac{1}{3} m g \sin \varphi$$

pogoj za »nepodrsavanje«: $F = \frac{1}{3} m g \sin \varphi < F_\ell = m g \cos \varphi k_\ell$

$$\tan \varphi < 3 k_\ell$$

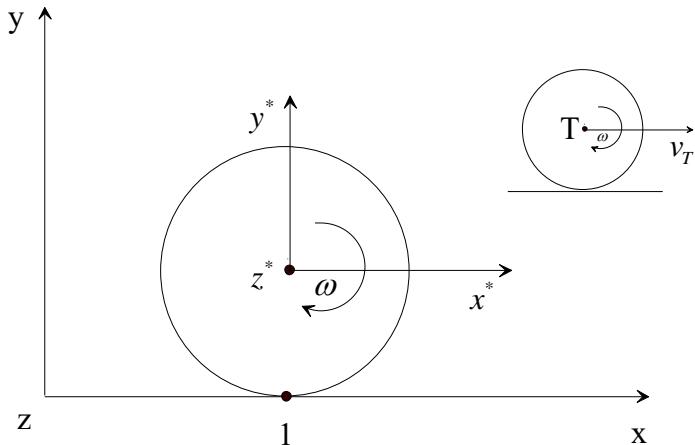
Poskus: kotaljenje po klancu: tekma valjev (poln in votel valj na strmini)

$$a_T = \frac{2}{3} g \sin \varphi$$

$$a_T = \frac{1}{2} g \sin \varphi$$



Dokaz veljavnosti enačbe: $v_T = R\omega$ (brez podrsavanja)



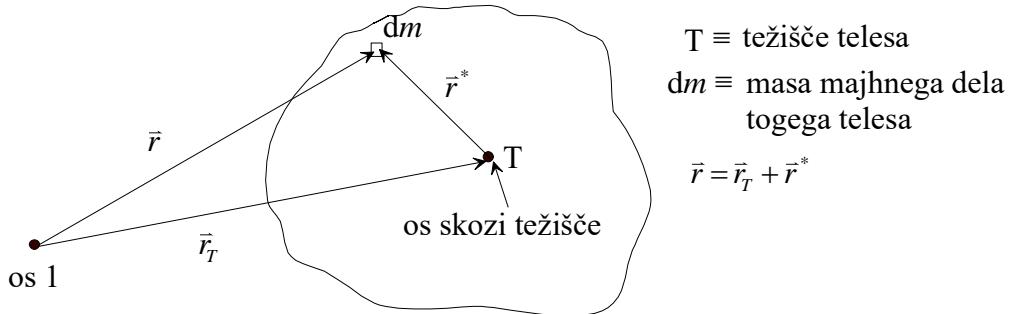
hitrost točke 1 v laboratorijskem sistemu: $v_1 = 0$

hitrost točke 1 merjena v težiščnem sistemu: $v_1^* = -R\omega$
 R = polmer preseka valja

Galilejeva transformacija za hitrosti: $v_1 = v_1^* + v_T$

$$v_T = -v_1^* = R\omega$$

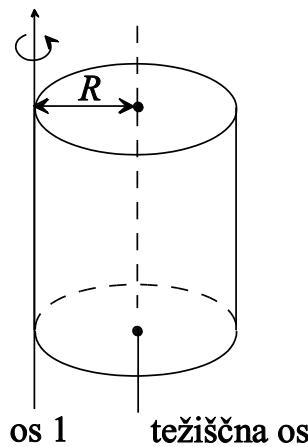
Steiner-jev izrek (povezuje vztrajnostni moment 2 togih teles okoli dveh paralelnih osi, od katerih poteka ena skozi težišče telesa)



$$\begin{aligned}
 J_1 &= \int r^2 dm = \int \vec{r} \cdot \vec{r} dm = \int (\vec{r}_T + \vec{r}^*) \cdot (\vec{r}_T + \vec{r}^*) dm = \\
 &= \int r_T^2 dm + 2 \vec{r}_T \cdot \int \vec{r}^* dm + \int r^{*2} dm = m r_T^2 + J^*
 \end{aligned}$$

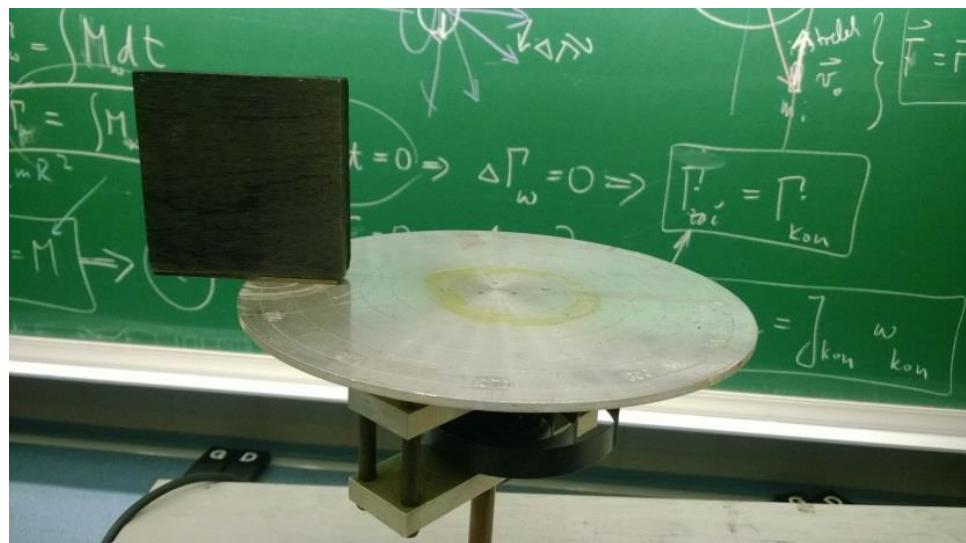
$$\vec{r}_T^* = \frac{1}{m} \int \vec{r}^* dm = 0$$

Primer:



$$J_1 = J^* + m r_T^2 = \frac{1}{2} m R^2 + m R^2 = \frac{3}{2} m R^2.$$

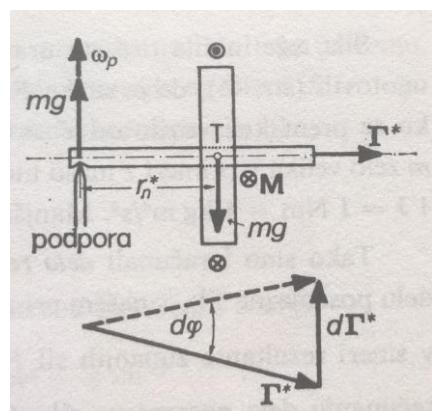
Poskus: streljanje v tarčo na mizici sučnega balističnega nihala



Poskus: mali bicikel



Poskus: precesija (vrtavka na vrvici)



Kinetična energija togega telesa v splošnem



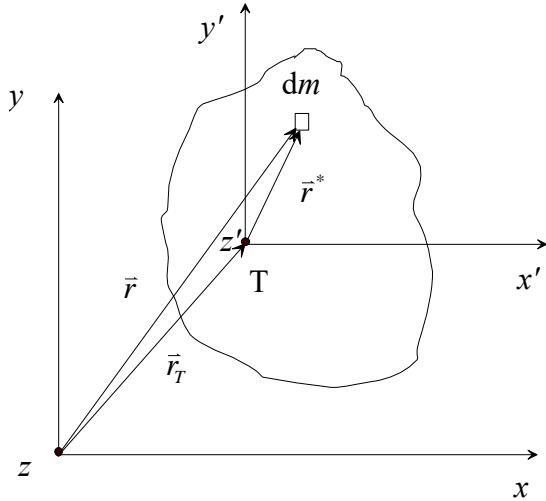
$$= \overline{v} || d\vec{r} | \cos \varphi = \overline{v} d\vec{v}$$

The diagram shows two semicircular paths on a chalkboard, representing the motion of a particle in a circle.



- izraz $\frac{1}{2}mv^2$ le, če se togo telo giblje po premici

predpostavimo, da so koordinatne osi laboratorijskega inercialnega koordinatnega sistema (x, y, z) paroma vzporedne z osmi lokalnega težiščnega sistema (x', y', z') :



$T \equiv$ težišče togega telesa

$$\vec{r} = \vec{r}_T + \vec{r}^*$$

$$W_k = \frac{1}{2} \int v^2 dm = \frac{1}{2} \int \vec{v} \cdot \vec{v} dm$$

$$W_k = \frac{1}{2} \int (\vec{v}_T + \vec{v}^*) \cdot (\vec{v}_T + \vec{v}^*) dm = \frac{1}{2} \int v_T^2 dm + v_T \cdot \int \vec{v}^* dm + \frac{1}{2} \int \vec{v}^{*2} dm$$

$$\vec{r}_T^* = \frac{1}{m} \int \vec{r}^* dm = 0 \quad \frac{d}{dt} \left(\frac{1}{m} \int \vec{r}^* dm \right) = \frac{1}{m} \int \vec{v}^* dm = 0$$

vrtenje okoli trenutne osi s kotno hitrostjo ω , ki je v izbranem trenutku enaka za vse točkaste mase dm : $v^* = R^* \omega$

$$W_k = \frac{1}{2} m v_T^2 + \frac{1}{2} J^* \omega^2$$

$$\frac{1}{2} \int v^{*2} dm = \frac{1}{2} \omega^2 \int R^{*2} dm = \frac{1}{2} J^* \omega^2 \quad J^* = \int R^{*2} dm$$

za izbrani težiščni sistem lahko v splošnem zapišemo rotacijsko kinetično energijo kot:

$$W_r = \frac{1}{2} \sum J_{\alpha\beta} \omega_\alpha \omega_\beta$$

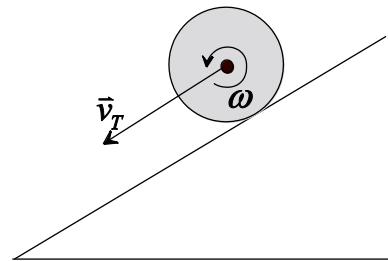
$$\alpha = x^*, y^*, z^*, \quad \beta = x^*, y^*, z^*$$

kjer so $J_{x^*x^*}, J_{y^*y^*}$ in $J_{z^*z^*}$ vztrajnostni momenti okrog treh pravokotnih osi x^*, y^* in z^* , $J_{\alpha\beta} (\alpha \neq \beta)$ pa so deviacijski momenti. Če so osi težiščnega sistema hkrati tudi glavne osi dobi izraz (29) obliko:

$$W_r = \frac{1}{2} J_{x^*} \omega_1^2 + \frac{1}{2} J_{y^*} \omega_2^2 + \frac{1}{2} J_{z^*} \omega_3^2$$

Primer : če se trenutna os vrtenja vedno ujema z eno izmed glavnih osi, n.pr. z J_{x^*}

$$W_r = \frac{1}{2} J_{x^*} \omega_1^2$$



$$v_T = R \omega \quad J = \frac{1}{2} m R^2$$

$$W_k = \frac{1}{2} m v_T^2 + \frac{1}{2} J \omega^2 = \frac{1}{2} m v_T^2 + \frac{1}{2} \left(\frac{1}{2} m R^2 \right) \frac{v_T^2}{R^2} = \frac{3}{4} m v_T^2$$

Zgled: vrtenje okoli ene izmed glavnih osi

če se trenutna os vrtenja vedno ujema z eno izmed glavnih osi :

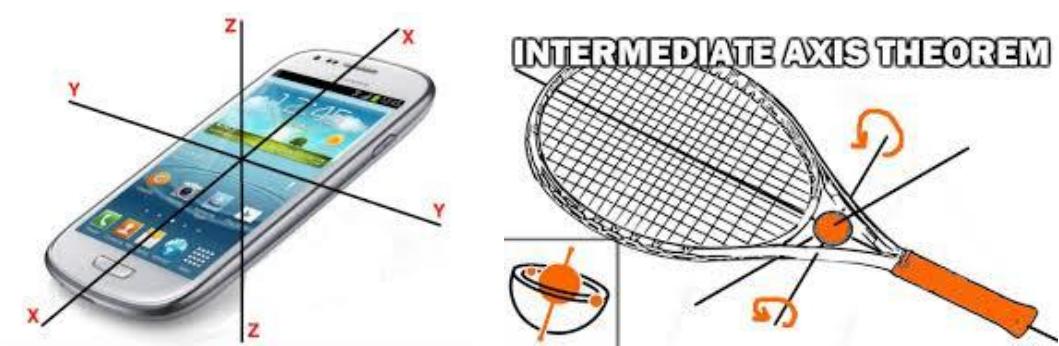
okoli x-osi: $W_r = \frac{1}{2} J_x \omega_1^2 = \frac{\Gamma^2}{2J_x}$

okoli y-osi: $W_r = \frac{1}{2} J_y \omega_2^2 = \frac{\Gamma^2}{2J_y}$

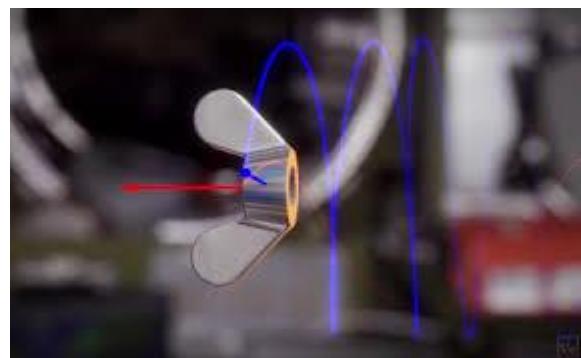
okoli z-osi: $W_r = \frac{1}{2} J_z \omega_3^2 = \frac{\Gamma^2}{2J_z}$



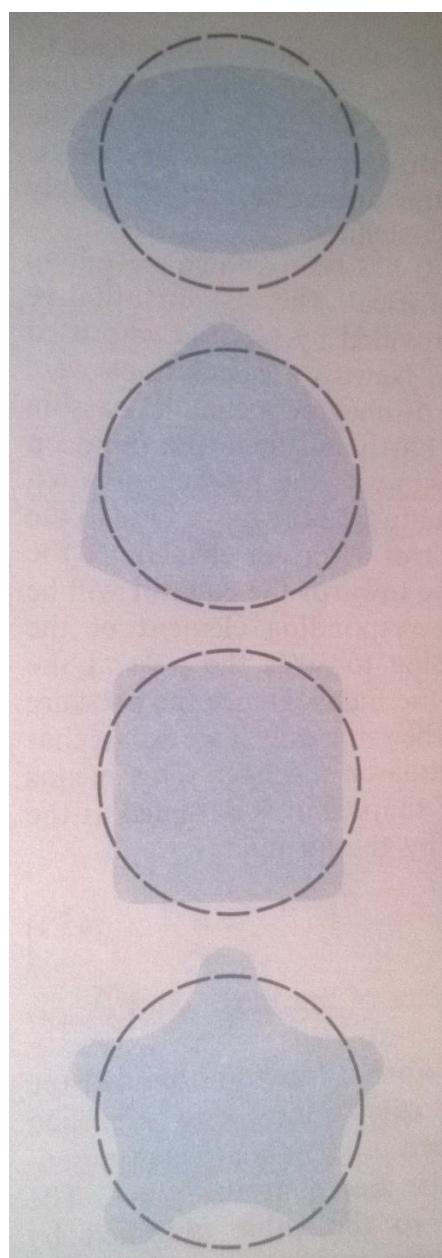
Dzhanibekov efekt



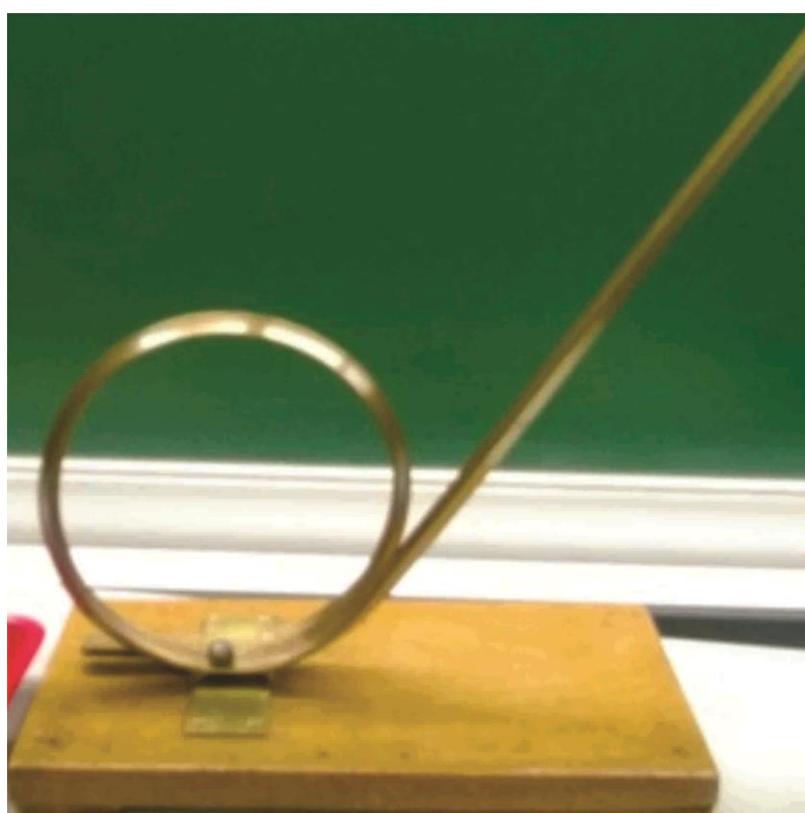
https://www.youtube.com/watch?v=1VPfZ_XzisU



oblika Zemlje :



ZGLED: »LOOPING« - zaviti žleb in kovinska kroglica, ohranitev vsote kinetične in gravitacijske potencialne energije



MEHANSKO RAVNOVESJE

$$\sum_i \vec{F}_i = 0 \quad \sum_i \vec{M}_i = 0$$

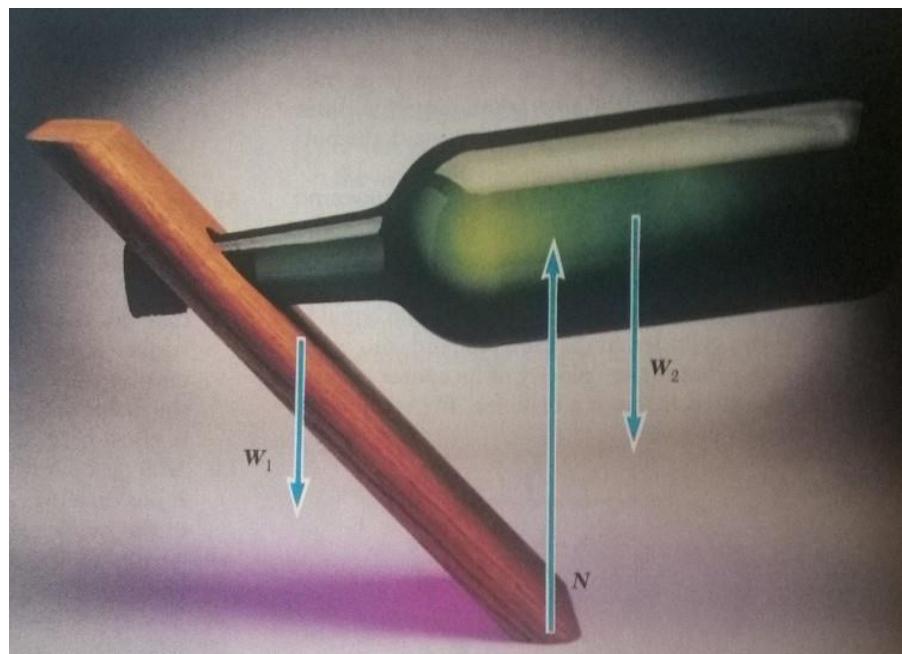


Foto: Serway